# Numerical Study of the Inclination Impact of an Inclined "T" Shaped Double Cavity on the Solution Symmetry and Heat Transfer 

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#### Abstract

This work has as principal goal the numerical investigation of mixed convective heat transfer in an inclined " T " shaped double cavity containing four heated blocks and four openings. The heated blocks are attached to the lower wall and maintained at a constant temperature $\theta_{H}$. The upper wall is maintained cold at a constant temperature $\theta_{C}\left(\theta_{\mathrm{C}}<\theta_{\mathrm{H}}\right)$. The vertical walls are rigid and adiabatic. The equations that rule the problem are solved by using the finite volume method and the SIMPLEC algorithm. We will put emphasis on the impact of the inclination angle on the solution symmetry and the heat transfer. The most important parameters are: the inclination angle $0^{\circ} \leq \varphi \leq 90^{\circ}$, the Rayleigh number $R a=2.10^{5}$, the Reynolds number $100 \leq \operatorname{Re} \leq 1000$, the Prandtl number $\operatorname{Pr}=0.72$, the height of the blocks $B=0.5$, the opening width $C=0.15$ and the distance between blocks $D=0.5$. Various results for flow field, isotherms and Nusselt number variations with $\operatorname{Re}$ and $\varphi$ are presented.


Keywords-Mixed Convection, inclination effect, "T" shaped double cavity, air jet

## I. INTRODUCTION

The failure rate of electronic devices ascends with temperature almost exponentially. Thermal cycling can cause break at connections. Therefore, thermal control has become increasingly important in the design and operation of electronic equipments in many fields of the new industry in order to reduce temperature and heat excess [1-3].

For that, many authors and scientists have put emphasis on convective heat transfer inside channels. Where, A.Korichi et al [4] have carried out numerically heat transfer and fluid flow in a horizontal channel containing two obstacles on the lower wall and one obstacle on the upper wall located in between the two. They found that the increase of the Reynolds number value, leads to the increase of the heat removed from the obstacles with a maximum heat removal around the obstacle corners. A numerical simulation [5] into an alternative method of natural convection enhancement by the transverse oscillations of a thin short plate, strategically positioned in
close proximity to a rectangular heat source was carried out. The results indicated that this method can serve as a feasible, simpler, more energy and space efficient alternative to common methods of cooling for low power dissipating devices operating at conditions just beyond the reach of pure natural convection.

However, a numerical study [6] of mixed convection heat transfer in an inclined rectangular channel with three heat sources on the lower wall was carried out using the finite element method and the Petrov-Galerkin technique. They obtained that the inclination has a stronger influence on the flow and heat transfer for low Reynolds numbers. In general, cases which show the lowest temperature distributions on the modules are those where the inclination angles are $45^{\circ}$ and $90^{\circ}$.

Among the studies in literature, there exist various investigations that treat heat transfer in cavities of different forms such as Rahman and al [7] that have analyzed numerically a computational work to show the effects of buoyancy ratio on heat and mass transfer in a triangular enclosure with zig-zag shaped bottom wall for different parameters. They obtained that average Nusselt and Sherwood numbers increase by $89.18 \%$ and $101.91 \%$ respectively as Br increases from -10 to 20 at $\mathrm{Ra}_{\mathrm{T}}=10^{6}$. Also, average Nusselt decreases by $16.22 \%$ and Sherwood numbers increases by $144.84 \%$ as Le increases from 0.1 to 20 at this Rayleigh number. Mahmoodi [8] has put emphasis on problem of natural convection fluid flow and heat transfer of Cu -water nanofluid inside L-shaped cavities. He has obtained that the mean Nusselt number for all rage of cavity aspect ratio increases with increase in the Rayleigh number and the solid volume fraction of the nanofluid. More, he has observed that the rate of heat transfer increases with decreasing the aspect ratio of the cavity.

Moreover, numerous researchers [9, 10, 11] have investigated convective heat transfer inside " T " form cavities. Meskini and al. [9] have examined numerically the
enhancement of heat transfer in a square cavity with identical heated rectangular blocks adjacent to its upper wall, and submitted to a vertical jet of fresh air from below. They found different solutions, when varying $\mathrm{Re}, \mathrm{Ra}$ and B , on which the resulting heat transfer, depends significantly.

Recently, various articles have appeared in literature, one of them the investigation [12] of mixed-convection fluid flow and heat transfer in a square cavity filled with $\mathrm{Al}_{2} \mathrm{O}_{3}$-water nanofluid using the finite volume method (FVM) and SIMPLER algorithm. They obtained that the average Nusselt number for all ranges of solid volume fraction increases with a decrease in the Richardson number. The results elucidate irregular changes of mean Nusselt number at different Richardson numbers versus variation of inclination angles in any case.

This paper has as purpose the numerical study of the inclination effect on mixed convection inside a " T " form double cavity which contains four openings and four heated blocks attached to its lower wall. The upper wall is maintained cold while the vertical walls and the lower wall are adiabatic as displayed in Fig. 1.

## II. Physical problem and mathematical formulation

The configuration shown in Figure 1 is considered. It is a "T" shaped double cavity containing four openings and four heated blocks arranged on its lower wall and it is maintained at a constant temperature $\theta_{\mathrm{H}}$. The upper wall is cold at $\theta_{\mathrm{C}}<\theta_{\mathrm{H}}$. The vertical walls and the lower wall are adiabatic. The height of the blocks is constant for $\mathrm{B}=0.5$.


Fig. 1 Schematic representation of the problem

## A. Mathematical Formulation

The governing equations of this problem are those of Navier-Stockes associated to the energy equation in unsteady state.

The fluid is supposed to be bidimensional, laminar and incompressible while Boussineq approximation is validated and the fluid properties are constant.
The local Nusselt numbers along the heated blocks $\mathrm{Nu}_{\mathrm{L}}, \mathrm{Nu}_{\mathrm{I}}$ and $\mathrm{Nu}_{\mathrm{R}}$ will be calculated for examining the heat exchange through the double cavity. Respectively, the mean Nusselt number over the active surfaces of the double cavity is given by:

$$
\begin{equation*}
N u=N u_{L}+N u_{I}+N u_{R} \tag{1}
\end{equation*}
$$

The global Nusselt number on the left block surface is:

$$
\begin{equation*}
N u_{L}=\left.\int_{0}^{B} \frac{\partial T}{\partial x}\right|_{x=0.25} d y+\left.\int_{0}^{0.25} \frac{\partial T}{\partial y}\right|_{y=B} d x \tag{2}
\end{equation*}
$$

The global Nusselt number on the intermediate block surfaces is:

$$
\begin{align*}
& N u_{I}=\left.\int_{0}^{B} \frac{\partial T}{\partial x}\right|_{x=0.75} d y+\left.\int_{0.75}^{1.25} \frac{\partial T}{\partial y}\right|_{y=B} d x  \tag{3}\\
& +\left.\int_{0}^{B} \frac{\partial T}{\partial x}\right|_{x=1.25} d y
\end{align*}
$$

The global Nusselt number on the right block surfaces is:

$$
\begin{equation*}
N u_{R}=\left.\int_{0}^{B} \frac{\partial T}{\partial x}\right|_{x=1.75} d y+\left.\int_{1.75}^{2} \frac{\partial T}{\partial x}\right|_{y=B} d x \tag{4}
\end{equation*}
$$

The dynamic and thermal boundary conditions associated to the problem equations are:

- $\theta=U=0, V=1$ for $\mathrm{y}=0 ; 0.425 \leq \mathrm{x} \leq 0.575$ and
$1.425 \leq x \leq 1.575$ (admission openings)
- $U=V=0, \frac{\partial \theta}{\partial y}=0$ for $\mathrm{y}=0 ; 0.25 \leq \mathrm{x} \leq 0.425 ; 0.575 \leq \mathrm{x} \leq 0.75$;
$1.25 \leq x \leq 1.425$ and $1.575 \leq x \leq 1.75$ (The parts of lower wall separating the opening and the blocks)
- $U=V=\frac{\partial \theta}{\partial x}=0$ for $\mathrm{B} \leq \mathrm{y} \leq 1 ; \quad \mathrm{x}=0$ and $\mathrm{x}=2 \quad$ (adiabatic vertical walls)
- $U=V=0$ on the rigid walls.

The $\theta, \mathrm{U}$ and V values in the evacuation opening are extrapoled by imposing the second derivates null with respect to $y$ as in $[13,14]$.

## B. Numerical Method

The finite volume method and the SIMPLEC algorithm will be used to solve the governing equations and to treat the Velocity-Pressure coupling [15,16] while, the schema Quick is chosen to discritisize the convective terms of the transport equation [17]. The grid size of $120 \times 81$ has been chosen to model accurately the flow fields. The time step considered in
this investigation is between $10^{-4}$ and $10^{-3}$. The code is validated by comparing our results with those reached by Kalache [18] in case of natural convection flows inside trapezoidal cavity and then they will be compared with results obtained in natural convection in a vertical channel by Destrayaud and Fichera [19]. It is worthy to say that we have obtained a good agreement in this comparison.

## III. RESULTS AND DISCUSSION

The results in terms of stream lines, isotherms and Nusselt number variations for inclination angle $0^{\circ} \leq \varphi \leq 90^{\circ}$, Reynolds number $100 \leq \operatorname{Re} \leq 1000$ and $\mathrm{Ra}=2.10^{5}$, will be displayed.

## A. Streamlines and isothermal lines

When we vary $\operatorname{Re}$ and $\varphi$, we remark the existence of different solutions: ICF (Intra Cellular Flow), ECF (Extra Cellular Flow) $[9,10]$ and a combination of FF (Forced Flow) and UCLF (Uni Cellular Left Flow) or ECF and UCLF [9]. More, we also remark that the solution of the double cavity is formed of two solutions, one in each sub cavity: SC1 (Simple Cavity 1) and SC2 (Simple Cavity 2). According to the flow field, it is observed that a communication is between the two adjacent cavities.


Fig. 2 Streamlines and isotherms for Rayleigh number $\mathrm{Ra}=2.10^{5}$; inclination angle $\varphi=0^{\circ}$ : \{(a),(b),(c) \}

In this section, we notice that all the solutions are symmetric, with respect to $P_{2}$ axis passing by the middle of the double cavity, for cases not inclined $\left(\varphi=0^{\circ}\right)$ as presented in Fig. 2. Where, at $\operatorname{Re}=200, \varphi=0^{\circ}$, Fig. 2a, we note that this solution is an ICF type and the solution symmetry is conserved with respect to $\mathrm{P}_{2}$ axis. The flow structure is composed of the Open lines (OL) that fill the whole double cavity surrounding two cells (in each domain SC 1 and SC 2 ). The corresponding thermal fields reveal that the horizontal surfaces of the heated blocks are badly ventilated due to the absence of the convective cells while, the vertical active ones are well ventilated because the air jet is still inside the two micro cavities. They also reveal that the isotherms have a
tooth shape inside each sub cavity. By increasing Reynolds number to $500, \varphi=0^{\circ}$, Fig. 2b, we perceive that the solution is an ICF type and symmetric. Fig. 2b presents the appearance of four Rayleigh-Bénard cells above the blocks. The open lines are parallel to $\mathrm{P}_{1}$ axis and pushed by the convective cells to the middle of SC1 and SC2. This figure also presents that the recirculation cells have increased in size. The corresponding isotherms show that the horizontal surfaces of the heated blocks are more cooled than the other ones because of the convective cells. This demonstrates that the symmetry of the solutions ICF is in favor of the cooling of the horizontal heated surfaces. The air jet arrives to the upper part of the double cavity.By more increasing Re to $1000, \varphi=0^{\circ}$, Fig. 2c, the solution (ICF) is seen symmetric with respect to $\mathrm{P}_{2}$ axis while the symmetry to $P_{1}$ axis is not conserved. The flow structure shows that the shape of the convective cells, which are above $B_{L}$ and $B_{R}$, has changed. The isothermal lines of this case indicate that the horizontal surfaces are well ventilated due to the rotation of the convective cells. The vertical active surfaces are badly ventilated because of the recirculation cells. They also indicate a sharp heat decrease of the temperature gradient above the block $\mathrm{B}_{\mathrm{I}}$. The flow is influenced by the convective and the recirculation cells.


Fig 3. Streamlines and isotherms for Rayleigh number $\mathrm{Ra}=2.10^{5}$; inclination angle $\varphi=36^{\circ}$ : $\{(\mathrm{a}),(\mathrm{b}),(\mathrm{c})\}$
When we incline the double cavity, $\varphi=36^{\circ}$, Fig. 3, we notice that the solution symmetry is destroyed. Where, at
$\operatorname{Re}=200, \varphi=36^{\circ}$, Fig. 3a, we remark that the solution is asymmetric (UCLF+ECF). The flow field is composed of a convective cell above $\mathrm{B}_{\mathrm{L}}$ and a small cell above $\mathrm{B}_{\mathrm{I}}$. The open lines are divided into two parts, one part fills (the middle of SC 1 ), the other one deflects in the corner of $\mathrm{B}_{\mathrm{I}}$ then turn around a cell for exiting from the evacuation opening (of SC 1 ). In addition, the OL fill the majority of SC2 domain, indicating a communication between the two sub cavities. The corresponding isotherms show a good ventilation of the vertical active walls. Most, the horizontal active wall of $B_{L}$ is more cooled than the other horizontal ones. For $\mathrm{Re}=500, \varphi=$ $36^{\circ}$, Fig .3b, we remark that the solution symmetry is not conserved. The type of this solution is ICF. The flow structure reveals that the right convective cell of SC1 has disappeared letting the space above $\mathrm{B}_{\mathrm{I}}$ to the convective cell (of SC2), which demonstrates a mutual communication between the two sub cavities. Besides, the flow in SC2 overtakes the flow in SC1 domain. It also reveals the Rayleigh Bénard cell above $B_{R}$ has decreased in size. The corresponding isotherms depict that the horizontal active surfaces of $\mathrm{B}_{\mathrm{L}}$ and $\mathrm{B}_{\mathrm{I}}$ are more ventilated than the horizontal surface of $B_{R}$. They also depict the bad ventilation of the vertical surfaces of the heated blocks due to the recirculation cells. By increasing Re to $1000, \varphi=$ $36^{\circ}$, Fig. 3c, we perceive that the solution symmetry is not conserved with respect to $\mathrm{P}_{2}$ axis and it is an ICF type. The type of this solution is conserved (ICF) while the form of the convective cells has changed. The isotherms of this case show that the horizontal active wall of $\mathrm{B}_{\mathrm{R}}$ is no more ventilated due to the small cell. The vertical active walls are badly ventilated.

$\mathrm{Re}=1000$
Fig . 4 Streamlines and isotherms for Rayleigh number $\mathrm{Ra}=2.10^{5}$; inclination angle $\varphi=72^{\circ}$ : \{(a),(b),(c) \}
By more inclining the double cavity $\varphi=72^{\circ}$, we observe that the solutions are still asymmetric. At $\operatorname{Re}=200, \varphi=72^{\circ}$, Fig. 4a, we remark that the solution is asymmetric with respect to $\mathrm{P}_{2}$ axis (UCLF+FF). The flow field exposes that the convective cells above $B_{L}$ and $B_{I}$ have increased in size. The recirculation cells have also increased in size and the small cell inside the OL of SC 2 has disappeared. The corresponding isotherms evince that the horizontal active surface of $B_{L}$ is more ventilated than the other horizontal one. The lower parts of the vertical walls of $\mathrm{B}_{\mathrm{L}}$ and $\mathrm{B}_{\mathrm{I}}(\mathrm{SC} 2)$ are badly ventilated. Most, the flow is influenced by the inclination. By increasing Re to $500, \varphi=72^{\circ}$, Fig. 4 b, we note that the solution is still asymmetric. The type of this solution is ICF. The flow structure indicates that the intensity of the convective cell above $B_{R}$ has increased. While, the intensity of the Rayleigh Bénard cell above $B_{I}$ is reduced letting space to a new cell to appear (in SC1). The right recirculation cell of SC2 has increased in size. The isotherms prove that the vertical walls are badly ventilated. They also prove the horizontal active surface of $B_{R}$ is less ventilated than the other horizontal ones. By increasing Re to $1000, \varphi=72^{\circ}$, Fig. 4c, we behold that the solution is not symmetric and it is an ICF type. The flow structure and the corresponding isotherms have not changed. They show the same phenomenon as the case of $\operatorname{Re}=1000, \varphi=$ $36^{\circ}$. It is worthy to say that the inclination angle has no effect on the solution symmetry for high Reynolds number Re and for $\varphi>0^{\circ}$.

## B. Heat transfer

We will present the variations of the global Nusselt numbers along the heated blocks $\mathrm{Nu}_{\mathrm{L}}, \mathrm{Nu}_{\mathrm{I}}$ and $\mathrm{Nu}_{\mathrm{R}}$, according to Reynolds number $100 \leq \operatorname{Re} \leq 1000$ and inclination angle $\varphi$ between $0^{\circ}$ and $90^{\circ}\left(\mathrm{Ra}=2.10^{5}\right)$, in Fig .5-7.

In one hand, we observe that the global Nusselt number along the left block $\mathrm{Nu}_{\mathrm{L}}$ falls gradually with Reynolds number Re , in the zone of mixed convection, and then it increases sharply for all inclination angles $36^{\circ} \leq \varphi \leq 90^{\circ}$ as mentioned in Fig .5. Most, it varies in the same manner for different inclination angle. However, the global Nusselt number along the right block $\mathrm{Nu}_{\mathrm{R}}$ goes up with Re , declines $\left(36^{\circ} \leq \varphi \leq 54^{\circ}\right)$ and then it increases, Fig .5. The global Nusselt number along the intermediate block $\mathrm{Nu}_{\mathrm{I}}$ rises steadily, Fig .6, decreases then it rockets with Re in the zone of the forced flow (in general, $500 \leq \operatorname{Re} \leq 1000$ ). Furthermore, the heat transfer is similar for different values of the inclination angle in the left, intermediate and right blocks.

In the other hand, Fig . 6 indicates that the global Nusselt number along the intermediate block $\mathrm{Nu}_{\mathrm{I}}$ varies differently with the inclination angle $\varphi$ for $100 \leq \operatorname{Re} \leq 500$. While it increases with $\varphi$ then it remains constant for $\mathrm{Re}=1000$ $\left(36^{\circ} \leq \varphi \leq 90^{\circ}\right)$. We note that there is a steady increase for the variation of the global Nusselt number along the left block $\mathrm{Nu}_{\mathrm{L}}$ with the inclination angle $\varphi$, Fig . 7 , for $300 \leq \operatorname{Re} \leq 1000$ in the mixed and forced convection zone.

Nonetheless, there is a slow rise then a gradual decrease with $\varphi$ for the lower Reynolds number $100 \leq \operatorname{Re} \leq 200$. The global Nusselt number along the right block $\mathrm{Nu}_{\mathrm{R}}$ diminish with the inclination angle $\varphi$ for $\mathrm{Re}=100$ while it goes up very slowly with $\varphi$ for $200 \leq \operatorname{Re} \leq 500$ as displayed in Fig .7. More, $\mathrm{Nu}_{\mathrm{R}}$ ascends steadily with $\varphi$ for high Reynolds number $\mathrm{Re}=700$.



Fig . $5 \mathrm{Nu}_{\mathrm{L}}$ and $\mathrm{Nu}_{\mathrm{R}}$ variations with Reynolds number Re and different values of inclination angle $\varphi\left(\mathrm{Ra}=2.10^{5}\right)$

## IV. CONCLUSION

This work simulates the influence of inclination angle on mixed convection inside a " T " shaped double cavity containing four heated blocks on its lower wall and four opening.
The results obtained, indicate:

- The existence of different solutions ICF, ECF and a combination of UCLF and FF or UCLF and ECF.
- The solution symmetry, with respect to $\mathrm{P}_{2}$ axis, is destroyed under the effect of inclination for cases with inclination.
- The symmetry of solutions ICF is in favor of the ventilation of the horizontal heated walls.
- When we increase Re, the inclination has no influence on the flow structure of the solution ( $\varphi>0^{\circ}$ ).


## Nomenclature

A Aspect ratio of calculation domain $A=L^{\prime} / H^{\prime}$
B Dimensionless block height $\left(\mathrm{B}=\mathrm{h}^{\prime} / \mathrm{H}^{\prime}\right)$
C Dimensionless openings diameter ( $\mathrm{C}=1^{\prime} / \mathrm{H}^{\prime}$ )
D Dimensionless space between blocks ( $\mathrm{D}=\mathrm{d}^{\prime} / \mathrm{H}^{\prime}$ )
d' Space between adjacent blocks (m)
H' Cavity height (m)
h' Blocks height, (m)
L’ Horizontal length of calculation domain (m)
l' Opening diameter (m)
$g \quad$ Acceleration due to gravity $\left(\mathrm{m} / \mathrm{s}^{2}\right)$
$\mathrm{Nu}_{\mathrm{L}} \quad$ Global Nusselt number along the planes of the left block
$\mathrm{Nu}_{\mathrm{I}} \quad$ Global Nusselt number along the planes of the intermediate block
$\mathrm{Nu}_{\mathrm{R}} \quad$ Global Nusselt number along the planes of the right block
$\theta \quad$ Dimensionless temperature, $\left(\mathrm{T}-\mathrm{T}_{\mathrm{C}}\right) /\left(\mathrm{T}_{\mathrm{H}}-\mathrm{T}_{\mathrm{C}}\right)$
T Temperature of the fluid
$\Delta \mathrm{T} \quad$ Temperature difference $\left(\mathrm{T}_{\mathrm{H}}-\mathrm{T}_{\mathrm{C}}\right)$
$\mathrm{U}_{0} \quad$ Characteristic velocity of the forced flow
$(\mathrm{x}, \mathrm{y})$ Dimensionless Cartesian coordinates in the two directions $(x, y)=\left(x^{\prime}, y^{\prime}\right) / H^{\prime}$
(u,v) Velocities
$(\mathrm{U}, \mathrm{V}) \quad$ Dimensionless velocities $(\mathrm{U}, \mathrm{V})=(\mathrm{u}, \mathrm{v}) / \mathrm{U}_{0}$

## Greek Symbols

$\mu$
$\rho \quad$ Volumetric coefficient of thermal expansion $\left(\mathrm{K}^{-1}\right)$
$\varphi \quad$ Inclination angle of the cavity
$\Lambda \quad$ Thermal conductivity $\left(\mathrm{Wm}^{1} \mathrm{~K}^{-1}\right)$
$v \quad$ Cinematic viscosity $\left(\mathrm{m}^{2} \mathrm{~s}^{-1}\right)$
P Fluid density $\left(\mathrm{kg} / \mathrm{m}^{3}\right)$
$\Psi \quad$ Dimensionless stream function, $\psi=\psi^{\prime} / \alpha$

## $\Omega \quad$ Dimensionless vorticity $\Omega=\Omega{ }^{\prime} \mathrm{H}^{\prime 2} / \alpha$

## Subscripts

H Heated wall
C Cold wall

## Non-dimensional Numbers

$\operatorname{Pr} \quad$ Prandtl number $(\operatorname{Pr}=v / \alpha)$
Re Reynolds number $\left(\operatorname{Re}=\mathrm{U}_{0} \times \mathrm{H}^{\prime} / v\right)$
$\mathrm{Ra} \quad$ Rayleigh number, $\left(\mathrm{Ra}=\mathrm{g} \beta \Delta \mathrm{TH}^{\prime} /(\alpha \nu)\right)$


Fig . $6 \mathrm{Nu}_{\mathrm{I}}$ variations with Reynolds number $\operatorname{Re}$ and the inclination angle $\varphi$ ( $\mathrm{Ra}=2.10^{5}$ )

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Fig . $7 \mathrm{Nu}_{\mathrm{L}}$ and $\mathrm{Nu}_{\mathrm{R}}$ variations with inclination angle $\varphi$ and different values of
 Reynolds number $\operatorname{Re}\left(\mathrm{Ra}=2.10^{5}\right.$

